**International Journal of Mathematics and Computer Research ISSN: 2320-7167 Volume 11 Issue 08 August 2023, Page no. 3694-3703 Index Copernicus ICV: 57.55, Impact Factor: 7.362**

**[DOI: 10.47191/ijmcr/v11i8.10](https://doi.org/10.47191/ijmcr/v11i8.010)**



# **Convective Cattaneo-Christov Heat Flux and Heat Generation Effect on Mhd Williamson Fluid Flow Over an Exponentially Stretching Surface with Thermal Radiation**

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**KEYWORDS:** Williamson fluid, Cattaneo-Christov heat flux, MHD, thermal radiation, HAM. **Mathematics Subject Classification:** 76

# **1 INTRODUCTION**

Williamson [1] proposed a fluid model known as Williamson fluid model to know the flow behaviour of pseudoplastic fluids. The heat transfer in viscoelastic flow caused by an exponentially stretched sheet is described by the Cattaneo-Christov heat flux model. Using the Oldroyds' upperconvected derivative, Christov proposed the frameindifferent generalisation of Cattaneo's law from which he derived a single temperature governing equation [2]. The upper-convected Maxwell model and Cattaneo–Christov heat flux model is used to investigate heat transfer and boundary layer flow of a viscoelastic fluid above a stretching plate with velocity slip boundary by Han et al. [3]. A modified version of Fourier's law known as the Cattaneo-Christov heat flow model, is used to investigate the phenomenon of heat transport by Khan et al.[4]. The two-dimensional Oldroyd-B fluid over a stretching surface with gyrotactic microorganism using Cattaneo-Christov heat flux is investigated by Bashir et al. [5]. Recently Bilal et al. [6] analyzed the two-dimensional incompressible flow of Williamson nanofluid over an exponentially stretching surface with Cattaneo- Christov heat flux model and reported that fluid velocity decreases in relation to the Williamson fluid model parameter.

The steady flow of nanofluid with gyrotactic microorganism over a Riga plate using Cattaneo- Christov

heat flux model is studied by Faizan et al. [7]. Alharbi et al. [8] used Tiwari-Das model and the Cattaneo-Christov model to examine the impact of Marangoni convection and volume fraction during heat transfer and divulged that Marangoni convection shortens skin friction. Mahabaleshwar et al. [9] compared the Cattaneo-Christov heat flux concept in the flow of two viscoelastic fluids to explore the heat transfer individualities with changing thermal conductivity using the characteristics of the Appell hypergeometric function. Shahzad et al. [10] explore the effect of gyrotactic microorganisms and convective thermal boundary conditions on the Darcy-Forchheimer in a micropolar nanofluid flow between two coaxial, parallel, and radially expanding disks and observed that the stretching ratio parameter of the disks accelerates the axial and micro rotational velocities of the nanofluid.

Reddy et al. [11] studied the two-dimensional viscous flow of Cassson nanofluid with Cattaneo–Christov model and used Spectral Homotopy Analysis Method (SHAM) to get the numerical solution. Ahmad et al. [12] investigated the Cattaneo-Christov heat flux model of second-grade nanofluid flow over a stretching sheet and used HAM to get analytic solution. Ahmed et al. [13] analyzed the Maxwell nanofluid flow in three-dimensional porous medium with the idea of Cattaneo-Christov and Buongiorno models

and the resultant non-linear ordinary differential equations are solved using Homotopy Analysis Method. Akinbo et al. [14] discussed the effect of viscous Walters B fluid flow of Cattaneo-Christov model over an exponentially stretching sheet and the results are compared using HAM and Galerkin Weighted Residual method. Amjad et al.[15] investigated the Cattaneo-Christov double diffusion (CCDD) heat flux model for Williamson nanofluid over an exponentially stretching surface with variable thermal conductivity.

Prasad et al. [16] observed that the distribution of temperature and concentration is decreased as a result of the Cattanneo-Christov heat flux model. The non-Newtonian fluid behavior of Casson fluid model over an Exponentially Stretching Sheet with Heat Source and Sink is studied by Prakash et al. [17]. Shehzad et al. [18, 19] analyzed Cattaneo-Christov model for both third-grade fluid flow over an exponentially stretching sheet and for Darcy-Forchheimer flow of an Oldroyd-B fluid over a moving sheet. The influence of Cattaneo-Christov model of Williamson fluid over a permeable sheet is studied using HAM by Ray [20]. The flow of Williamson Sutterby nanofluid in Darcy– Forchheimer sponge medium with Cattaneo-Christov heat flux is discussed by Yahya et al. [21]. The flow of Williamson and Casson fluid flows over a penetrable extending Sheet in a permeable medium using Runge-Kutta Fehlberg method together with Shooting method is studied by Mangathai et al. [22].

Jamshed et al. [23] applied Keller-box method in the analysis of engine oil-based Williamson hybrid nanofluid flow. Hayat et al. [24] surveyed the flow of thixotropic fluid over a stretching surface of Cattaneo-Christov heat flux model. The bioconvection nanofluid flow with temperaturedependent variable viscosity by applying HAM is studied by Mondal et al. [25]. The MHD three-dimensional flow of Maxwell fluid over a bi-directional stretching surface with Cattaneo-Christov heat flux model by using HAM is analyzed by Rubab et al. [26]. Salahuddin et al. [27] analyzed the thermal relaxation time effect of MHD Williamson fluid flow using Cattaneeo-Christov heat flux model over a stretching sheet. Kumar et al. [28] utilized Cattaneo-Christov model for the flow over a wedge and a cone and the results are compared using R-K method and Newton's method. Bhatti et al. [29] discussed the flow of Williamson nanofluid under the influence of thermal diffusion and thermal radiation over a porous stretching and shrinking sheet and the results are obtained using Successive linearization method (SLM) and Chebyshev spectral collocation method (CSC).

In this paper, we are analyzing Williamson fluid model in a boundary layer over an exponentially stretching surface in the presence of magnetic field, heat source and thermal radiation. The distribution of the paper is first section contains Introduction, second section contains Mathematical formulation, third section is Homotopy Analysis solution of the problem, fourth section is result and discussion and fifth

section consists of graphs and tables.

#### **2 MATHEMATICAL FORMULATION**

The two-dimensional steady heat and mass transfer of incompressible Williamson fluid flow of uniform magnetic field over an exponentially stretching sheet is considered. By considering  $u$  and  $v$  as velocity components along  $x$  and  $y$ directions and the fundamental equations for this model will be [4].

$$
\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0,\tag{1}
$$

$$
u\frac{\partial u}{\partial x} + v\frac{\partial u}{\partial y} = v\frac{\partial^2 u}{\partial y^2} + \sqrt{2}v\Gamma\frac{\partial u}{\partial y}\frac{\partial^2 u}{\partial y^2} - \frac{\sigma B_0^2}{\rho}u + g\beta(T - T_\infty) + g\beta^*(C - C_\infty),
$$
\n(2)

$$
u\frac{\partial T}{\partial x} + v\frac{\partial T}{\partial y} = \frac{k}{\rho c_p}\frac{\partial^2 T}{\partial y^2} - \lambda_1 \left[ u^2 \left( \frac{\partial^2 T}{\partial x^2} \right) + 2uv \left( \frac{\partial^2 T}{\partial x \partial y} \right) + v^2 \left( \frac{\partial^2 T}{\partial y^2} \right) \right] - \lambda_1 \left[ \left( u\frac{\partial u}{\partial x} + v\frac{\partial u}{\partial y} \right) \left( \frac{\partial T}{\partial x} \right) + \left( u\frac{\partial v}{\partial x} + v\frac{\partial v}{\partial y} \right) \left( \frac{\partial T}{\partial y} \right) \right] + \frac{\rho_0}{\rho c_p} (T - T_\infty) + \frac{\nu}{c_p} \left( \frac{\partial u}{\partial y} \right)^2 - \frac{1}{\rho c_p} \left( \frac{\partial q_r}{\partial y} \right), \tag{3}
$$

$$
u\frac{\partial c}{\partial x} + v\frac{\partial c}{\partial y} = D_m \frac{\partial^2 c}{\partial y^2} + D_T \frac{\partial^2 T}{\partial y^2},
$$
  
with boundary conditions, (4)

 $u = U_w(x) = U_0 exp\left(\frac{x}{l}\right)$  $\left(\frac{x}{l}\right)$ ,  $v=0$ ,  $-k\frac{\partial T}{\partial y}=h_f(T_w-T)$ ,  $C=$  $C_w$  at  $y = 0$ ;

$$
u \to 0, T \to T_{\infty}, C \to C_{\infty} \text{ as } y \to \infty.
$$
 (5)

Where  $\nu$ , the kinematic viscosity,  $k$ , thermal conductivity, Γ, time rate constant,  $ρ$ , density of the fluid,  $σ$ , the electrical conductivity,  $g$ , gravitational force,  $B_0$ , uniform magnetic field,  $\beta$ , thermal expansion coefficient,  $\beta^*$ , concentration expansion coefficient,  $T$ , the temperature of the fluid,  $C$ , the concentration of the fluid,  $c_p$ , specific heat,  $\lambda_1$ , thermal relaxation time,  $Q_0$ , heat source,  $D_m$ , mass diffusion coefficient,  $D_T$ , thermal diffusion coefficient,  $h_f$ , convective heat transfer coefficient,  $l$ , characteristic length, radiative heat flux  $q_r = -\frac{4\sigma^*}{3k^*}$ 3 ∗  $\frac{\partial T^4}{\partial y}$ , where  $\sigma^*$ , Stefan Boltzmann constant and  $k^*$ , absorption coefficient. When temperature difference is small,  $q_r$  can be linearized by expanding  $T^4$ into Taylor's series about  $T_{\infty}$ , takes the form after neglecting higher order terms by  $T^4 \cong 4T_{\infty}^3T - 3T_{\infty}^4$ .

Here we consider similarity transformations as follows,

$$
\eta = y \sqrt{\frac{U_0}{2\nu l}} exp\left(\frac{x}{2l}\right), u = U_0 exp\left(\frac{x}{l}\right) f'(\eta), v
$$

$$
= -\sqrt{\frac{\nu U_0}{2l}} exp\left(\frac{x}{2l}\right) [f(\eta) + \eta f'(\eta)],
$$

$$
\theta(\eta) = \frac{T - T_{\infty}}{T_{w} - T_{\infty}}, \phi(\eta) = \frac{C - C_{\infty}}{C_{w} - C_{\infty}}.
$$
\n
$$
(6)
$$

Equations (2) - (5) will get reduced to  $(7)$  - (10) by using (6)  $f''' + ff'' - 2(f')^2 + \lambda f''f''' - Mf' + Gr\theta +$  $Gm\phi = 0$ ,

$$
\theta'' - \frac{1}{2} Pr \gamma f^2 \theta'' + R \theta'' + \frac{1}{2} Pr \gamma f f' \theta' + Pr f \theta' + 2Q Pr \theta + Pr E c(f'')^2 = 0,
$$
\n(8)

$$
\phi'' + Scf \phi' + SoSc\theta'' = 0.
$$
\n(9)

at  $\eta = 0$ ,

$$
f'\to 0, \theta\to 0, \phi\to 0, \text{ as } \eta\to
$$

 $f = 0, f' = 1, \theta' = -\gamma_B [1 - \theta], \phi = 1,$ 

 $\infty$ . (10)

Where  $M = \frac{2 \sigma B_0^2 l}{r^2}$  $\frac{\partial \sigma B_0^2 l}{\partial v_w}$ ; Magnetic field parameter,  $\lambda = \Gamma \frac{v_w^2}{\sqrt{vl_w^2}}$ 3  $\frac{v_w}{\sqrt{vl}}$ ; Williamson fluid parameter,  $Pr = \frac{v \rho c_p}{h}$  $\frac{\partial c_p}{\partial k}$ ; Prandtl number,  $Gr = \frac{2g\beta(T_W - T_{\infty})l}{H^2}$  $\frac{w(T_w-T_{\infty})l}{U_w^2}$ ; Grashof number,  $Gm = \frac{2g\beta^*(c_w-c_{\infty})l}{U_w^2}$  $\frac{C_W-C_\infty}{U_W^2}$ ; modified Grashof number,  $\gamma = \frac{\lambda_1 U_w}{I_w}$  $\frac{\partial w}{\partial t}$ ; Thermal relaxation time parameter,  $\gamma_B = \frac{h_f}{k}$  $\frac{\partial f}{\partial k}$   $\sqrt{\frac{2\nu l}{U_0}}$  $\frac{2vl}{U_0}$ ; Biot number,  $R = \frac{16}{3}$ 3  $\sigma^* T^3_\infty$  $\frac{1}{kk^*}$ ; Radiation parameter,  $Q = \frac{Q_0}{Q_0}$  $\frac{Q_0}{\rho c_p U_w}$ ; Heat source parameter  $Ec = \frac{v_w^2}{\sqrt{v_w^2}}$  $\frac{U_W^2}{c_p(T_W - T_\infty)}$ ; Eckert number,  $Sc = \frac{v}{D\tau}$  $\frac{v}{Dm}$ ; Schmidt number,  $\mathcal{S}o = \frac{Dt(T_w - T_{\infty})}{w(G - G_v)}$  $\frac{\partial u(T_W - T_{\infty})}{\partial v(T_W - T_{\infty})}$ ; Soret number.

### **3. HOMOTOPY ANALYSIS SOLUTION**

Shijun Liao (1992) [31, 32, 33, 34, 35, 36, 37, 38, 39, 40] explained Homotopy Analysis Method (HAM) to solve nonlinear differential equations analytically. Using this method [41, 42] we solve coupled nonlinear equations of this problem. The steps of the method are,

$$
N[f(\eta)] = f''' + ff'' - 2(f')^{2} + \lambda f''f''' - Mf' + Gr\theta + Gm\phi,
$$
\n(11)

$$
N[\theta(\eta)] = \theta'' - \frac{1}{2} \text{Pr} \gamma f^2 \theta'' + R\theta'' + \frac{1}{2} \text{Pr} \gamma f f' \theta' +
$$
  
Prf\theta' + 2Q \text{Pr}\theta + \text{Pr} \text{Ec}(f'')^2, (12)

$$
N[\phi(\eta)] = \phi'' + Scf\phi' + SoSc\theta''.
$$
 (13)

Linear operators considered are as follows,

$$
L(f) = \frac{\partial^3 f}{\partial \eta^3} + \frac{\partial^2 f}{\partial \eta^2},\tag{14}
$$

$$
L(\theta) = \frac{\partial^2 \theta}{\partial \eta^2} + \frac{\partial \theta}{\partial \eta'},
$$
\n(15)

$$
L(\phi) = \frac{\partial^2 \phi}{\partial \eta^2} + \frac{\partial \phi}{\partial \eta'},
$$
 (16)

which gives initial approximations as,

$$
f_0 = 1 - e^{-\eta}, \tag{17}
$$

$$
\theta_0 = \left(\frac{\gamma_B}{1 + \gamma_B}\right) e^{-\eta},\tag{18}
$$

$$
\phi_0 = e^{-\eta}.\tag{19}
$$

The nonlinear equations for approximate solutions are,

$$
(1-p)L[f(\eta, p) - f_0(\eta)] =
$$
\n
$$
hp\left[\frac{\partial^3 f}{\partial \eta^3} + f\frac{\partial^2 f}{\partial \eta^2} - 2\left(\frac{\partial f}{\partial \eta}\right)^2 + \lambda \frac{\partial^2 f}{\partial \eta^2}\frac{\partial^3 f}{\partial \eta^3} - M\frac{\partial f}{\partial \eta} + Gr\theta +
$$
\n
$$
Gm\phi\right],
$$
\n
$$
(1-p)L[\theta(\eta, p) - \theta_0(\eta)] = hp\left[\frac{\partial^2 \theta}{\partial \eta^2} - \frac{1}{2}Pr\gamma f^2\frac{\partial^2 \theta}{\partial \eta^2} +
$$
\n
$$
R\frac{\partial^2 \theta}{\partial \eta^2} + \frac{1}{2}Pr\gamma f\frac{\partial f}{\partial \eta}\frac{\partial \theta}{\partial \eta} +
$$
\n
$$
Prf\frac{\partial \theta}{\partial \eta} + 2QPr\theta + PrEc\left(\frac{\partial^2 f}{\partial \eta^2}\right)^2],
$$
\n
$$
(1-p)L[\phi(\eta, p) - \phi_0(\eta)] = hp\left[\frac{\partial^2 \phi}{\partial \eta^2} + Scf\frac{\partial \phi}{\partial \eta} +
$$

$$
Sosc\frac{\partial^2 \theta}{\partial \eta^2},\tag{22}
$$

with following boundary conditions,

$$
f(0, p) = 0, f_{\eta}(0, p) = 1, f_{\eta}(\infty, p) = 0,
$$
 (23)

$$
\theta_{\eta}(0,p) = -\gamma_B [1 - \theta(0,p)], \theta(\infty, p) = 0, \qquad (24)
$$

$$
\phi(0, p) = 1, \phi(\infty, p) = 0. \tag{25}
$$

Varying the values of  $p$  from 0 to 1 we get the solution from first approximation to required solution. Using Maclaurin's series expansion and applying Leibnitz theorem we get the series solution. The convergence of the series solution is derived by calculating the convergence parameter ℎ.

$$
L[f_m - \chi_m f_{m-1}] = hr_m(\eta),\tag{26}
$$

$$
L[\theta_m - \chi_m \theta_{m-1}] = h s_m(\eta), \tag{27}
$$

$$
L[\phi_m - \chi_m \phi_{m-1}] = ht_m(\eta), \tag{28}
$$

where 
$$
\chi_m = \begin{cases} 0, & \text{when } m \le 1 \\ 1, & \text{when } m > 1 \end{cases}
$$
 and (29)

$$
r_m(\eta) = f_{m-1}'''(\eta) + \sum_{k=0}^{m-1} f_{m-1-k}(\eta) f_k''(\eta)
$$
  
- 2 
$$
\sum_{k=0}^{m-1} f_{m-1-k}'(\eta) f_k'(\eta) - M f_{m-1}'(\eta)
$$
  
+  $\lambda \sum_{k=0}^{m-1} f_{m-1-k}''(\eta) f_k'''(\eta) + G r \theta_{m-1}(\eta) +$ 

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$$
Gm\phi_{m-1}(\eta),\tag{30}
$$

$$
s_m(\eta)
$$

$$
= \theta_{m-1}''(\eta) - \frac{1}{2} Pr\gamma \sum_{k=0}^{m-1} f_{m-1-k}(\eta) \sum_{j=0}^{k} f_{k-j}(\eta) \theta_{j}''(\eta) + R\theta_{m-1}''(\eta) + \frac{1}{2} Pr\gamma \sum_{k=0}^{m-1} f_{m-1-k}(\eta) \sum_{j=0}^{k} f'_{k-j}(\eta) \theta_{j}'(\eta) + PrEc \sum_{k=0}^{m-1} f_{m-1-k}(\eta) \theta_{k}'(\eta) + 2QPr\theta_{m-1}(\eta) +
$$
  
\n
$$
PrEc \sum_{k=0}^{m-1} f'_{m-1-k}(\eta) f'_{k}''(\eta), \qquad (31)
$$

 $t_m(\eta) = \phi_{m-1}''(\eta) + Sc \sum_{k=0}^{m-1} f_{m-1-k}(\eta) \phi_k'(\eta) +$  $Sosc \theta''_{m-1}$  $(\eta)$ , (32)

with boundary conditions,  

$$
f_m(0) = 0, f'_m(0) = 0, f'_m(\infty) = 0,
$$
 (33)

$$
\theta'_{m}(0) - \gamma_{B}\theta_{m}(0) = 0, \theta_{m}(\infty) = 0, \qquad (34)
$$

$$
\phi_m(0) = 0, \phi_m(\infty) = 0. \tag{35}
$$

The required solution is

$$
f = f_0 + f_1 + f_2 + \dots,
$$
 (36)

$$
\theta = \theta_0 + \theta_1 + \theta_2 + \dots,\tag{37}
$$

$$
\phi = \phi_0 + \phi_1 + \phi_2 + \dots \tag{38}
$$

Soving equtions (26)-(28) by using MATHEMATICA we get  $f_1 =$  $-\frac{-h+4Gmh-4hM-h\lambda-h\gamma_B+4Gmh\gamma_B+4Grh\gamma_B-4hM\gamma_B-h\lambda\gamma_B}{4(1+\lambda)}$  $\frac{4(1+\gamma_B)}{4(1+\gamma_B)} +$ 1  $\frac{1}{4(1+\gamma_B)}(e^{-2\eta})(-2e^{\eta}(h+2Gmh-2hM+h\lambda+$  $h\gamma_B + 2Gmh\gamma_B + 2Grh\gamma_B - 2hM\gamma_B + h\lambda\gamma_B$ ) +  $h((1 +$  $\lambda$ )(1 +  $\gamma_B$ ) + 4e<sup>n</sup>(Gr $\gamma_B$  + Gm(1 +  $\gamma_B$ ) – M(1 +  $\gamma_B$ ))(2 +  $(\eta))$ ),...  $\theta_1 = \frac{1}{4(1+1)}$  $\frac{1}{4(1+\gamma_B)}(e^{-2\eta})(hPr(2+\gamma)\gamma_B + 2EchPr(1+\gamma_B) 2e^{\eta}(\frac{1}{2})$  $\frac{1}{2}$ (2EchPr – 4h $\gamma_B$  + 6hPr $\gamma_B$  + 2EchPr $\gamma_B$  –  $8hPrQ_{Y_B} - 4hR_{Y_B} + 3hPr(Y_B) + h(2(1 + R) + Pr(-2 +$  $(4Q - \gamma)y_B(1 + \eta))$ ,... (40)  $\phi_1 = \frac{1}{2(1+1)}$  $\frac{1}{2(1+\gamma_B)}(e^{-2\eta})(hSc(1+\gamma_B)-2e^{\eta}(\frac{1}{2}))$  $\frac{1}{2}(-2h+3hSc 2h\gamma_B + 3hSc\gamma_B - 2hSoSc\gamma_B$ ) +  $h(1 + \gamma_B + Sc(-1 + (-1 +$  $(So)\gamma_B$ ))(1 +  $\eta$ ))),... (41)

# **4. RESULTS AND DISCUSSION**

The semi-analytical solutions of non-dimensional equations (7)-(9) are obtained by homotopy Analysis Method and numerical solution is obtained by Runge-Kutta method. The calculations were done using MATHEMATICA to derive the Williamson fluid's flow, heat, and mass transfer characteristics. The auxiliary non-zero parameters  $h_f$ ,  $h_\theta$ and  $h_{\phi}$  have significant influence on convergence of HAM solution. In figure 1 it can be seen convergence range for velocity profile is  $-1.5 < h_f < 0.5$ , for temperature profile is  $-1.5 < h_\theta < 1.2$  and for concentration profile is  $-2.5 <$  $h_{\phi}$  < 0.5.

Figures 2 - 13 delineates the effect of various parameters like Magnetic parameter  $M$ , Williamson parameter  $\lambda$ , Grashof number  $Gr$ , modified Grashof number  $Gm$ , Prandtl number Pr, Thermal relaxation time parameter  $\gamma$ , Heat source parameter  $Q$ , Biot number  $\gamma_B$ , Radiation parameter R, Eckert number  $Ec$ , Schmidt number  $Sc$  and Soret number So on velocity, temperature and concentration profiles. Figure 2 illustrates that the velocity of Williamson fluid drastically decreases as the effect of magnetic field increases. This is due to the Lorentz force which is dragging force in the presence of magnetic field. A similar pattern is observed in the case of Williamson parameter due to fluid relaxation time, shown in figure 3. From figures 4 and 5 we have observed that velocity profile increases with increase in Grashof number and modified Grashof number.

From figure 6, it is observed temperature profile decreases with increase in Prandtl number. The effect of Thermal relaxation parameter on the temperature field is shown in figure 7 and observed that temperature is seen to fall as a result of the relaxation parameter. In the absence of Thermal relaxation parameter, the Cattaneo-Christov heat flux model can be reduced to the Fourier's law of heat conduction. Additionally, Cattaneo-Christov heat flux model has a lower temperature than the Fourier's model. The temperature profile is clearly a function of the Heat source parameter in an increasing manner shown in figure 8. It is observed that the increase in values of Biot number, Radiation parameter and Eckert number increases the temperature profile seen in figures 9, 10 and 11.

The concentration profile increases with the increase in Williamson parameter and Soret number shown in figures 12 and 14, where as decreases with increase in Schmidt number and modified Grashof number shown in figures 13 and 15. Figures 16, 17 and 18 presents Domb-Sykes plots of velocity, temperature and concentration through which we get the radius of convergence as 0.71145, 0.12205 and 1.78374 respectively.

We have compared HAM solution with numerical solution for velocity, temperature and concentration profiles shown in figures 19, 20 and 21 and observed good agreement. We have also compared our solutions with the well known results of Elbashbeshy [30] shown in the table 1.

### **5. GRAPHS AND TABLES**



**Figure 4: Effect of Grashof number** *Gr* **Figure 5: Effect of modified Grashof number on**  $\blacksquare$ **velocity profile** *on velocity profile on velocity profile* 

'n

 $0.1$ 

 $\circ$ 

 $0.2$ 

 $0.0$ 

 $\mathbf{1}$ 

 $0.8$ 

 $\alpha$  $\mathbf{f}^{\star}\left( \eta\right)$  $\ddot{\phantom{0}}$ 

 $\overline{0}$ 

 $0.0$ 

 $\mathbf{f}^{\prime}(\eta)$  $\alpha$ 



**Figure 10:** Effect of Radiation parameter *R* Figure 11: Effect of Eckert number *Ec* on on temperature profile **temperature profile**







**Figure 18: Domb-Sykes concentration plot Figure 19: Comparison of HAM and numerical solution for Magnetic parameter**  $M = 0$ 



**Figure 19: Comparison of HAM and numerical Figure 20: Comparison of HAM and numerical solution for Magnetic parameter**  $M = 0$  solution for Magnetic parameter  $Sc = 1$ **on temperature profile on concentration profile**

**Table 1: Comparision of**  $-f''(0)$  **when**  $\lambda = M = Gr = Gm = 0$  **and**  $h = -0.5$ 

f''(0)		
Elbashbeshy [30]	Present work (HAM)	Present work (R K Method)
1.28181	.28182	1.28182

#### **Conflict of interest**

The authors declare that there is no conflict of interest.

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